

Reconstructing quantum states efficiently

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on work with
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- Full quantum state tomography is essential for present and future quantum devices

Quantum simulator: Preparation of elaborate states and operations on them

- Has the intended state indeed been prepared?
- Do the operations do what they are supposed to?

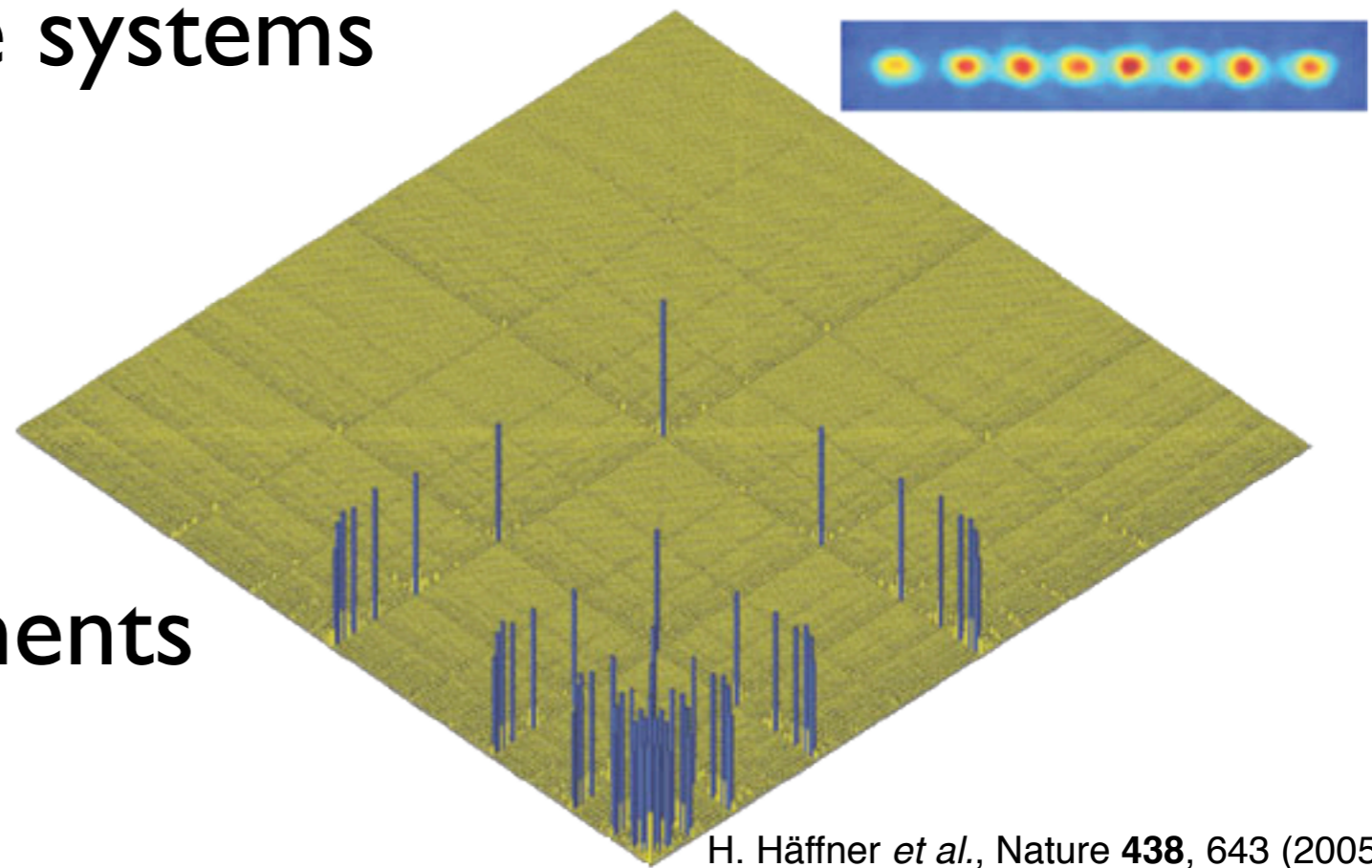
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- Possible for few-particle systems

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Measure complete

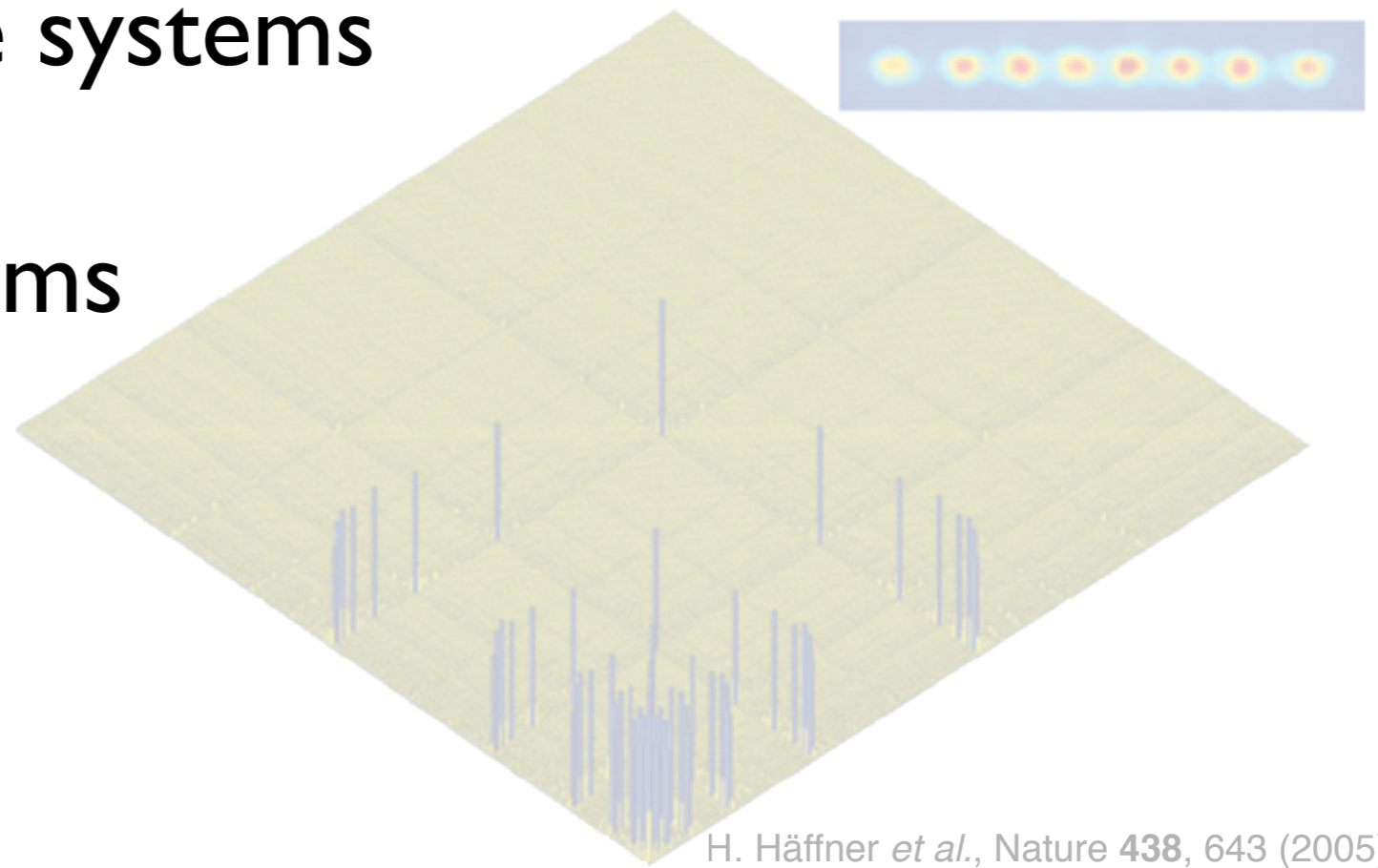
basis $\hat{P}_k = \bigotimes_{i=1}^N \hat{\sigma}_i^{\alpha_i}$

→ 4^N measurements

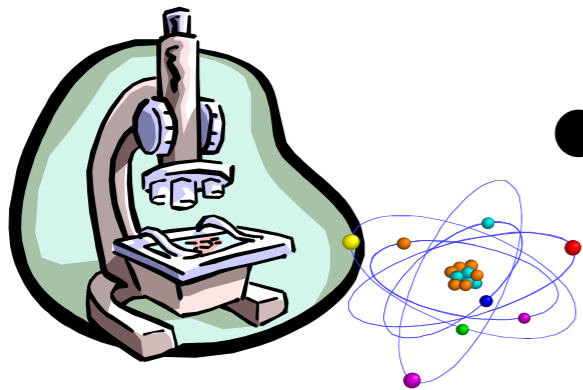


H. Häffner *et al.*, Nature **438**, 643 (2005).

- Full quantum state tomography is essential for present and future quantum devices
- Possible for few-particle systems
- Infeasible for large systems



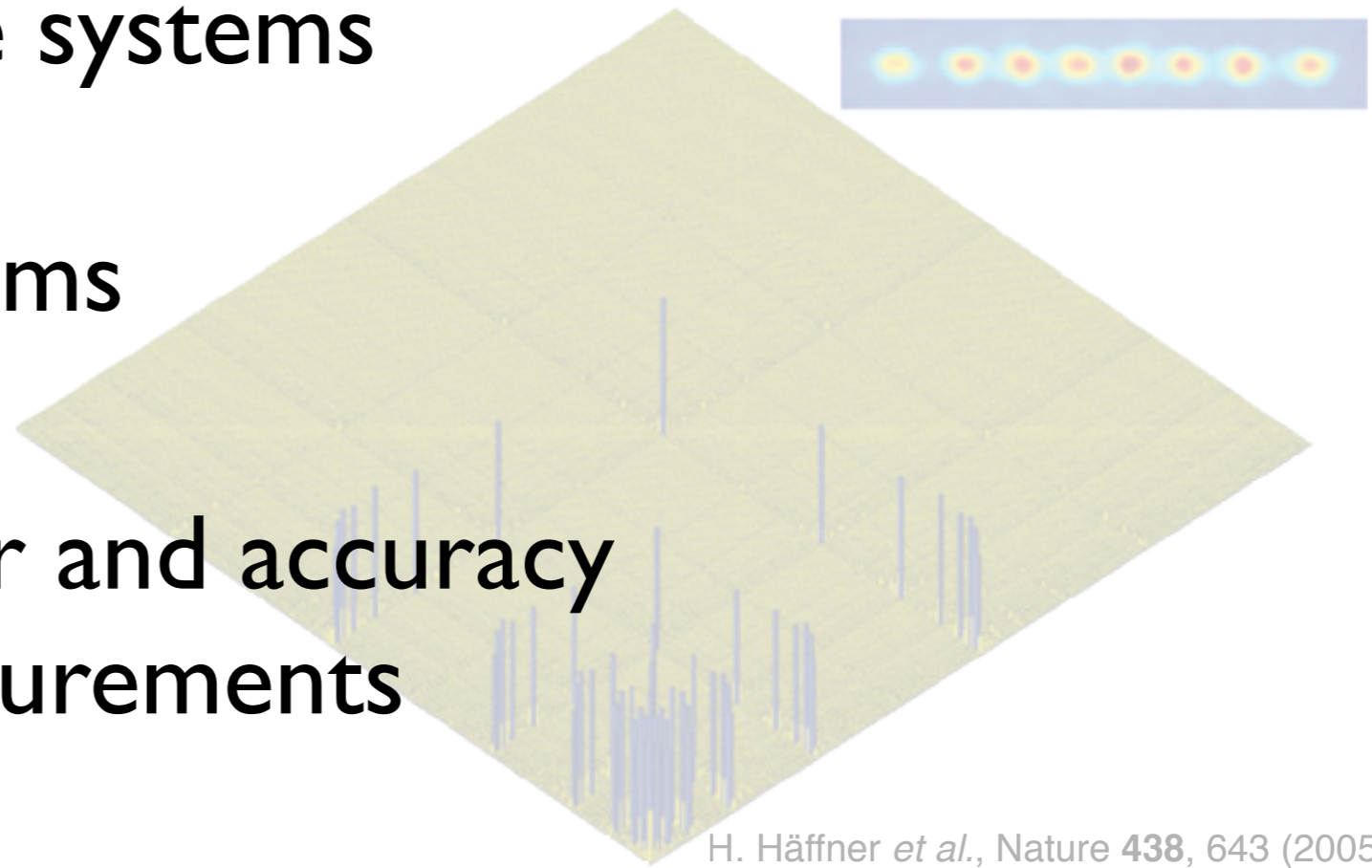
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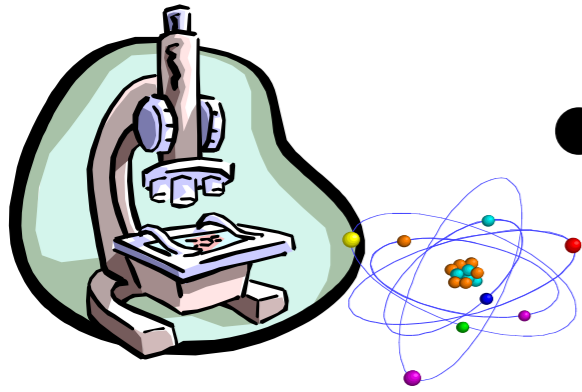
- Number and accuracy of measurements



- Find compatible state
- Storage space



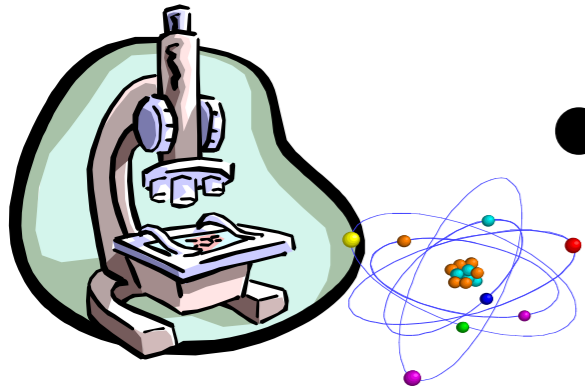
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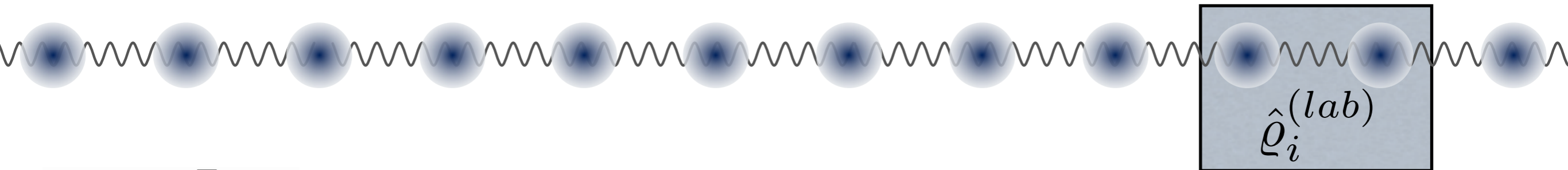
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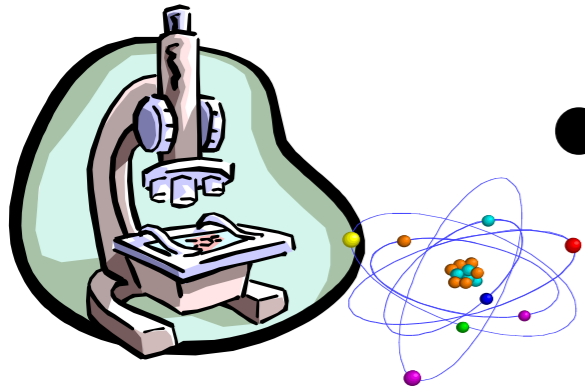
- Number and accuracy of measurements



Take only $\sim N$



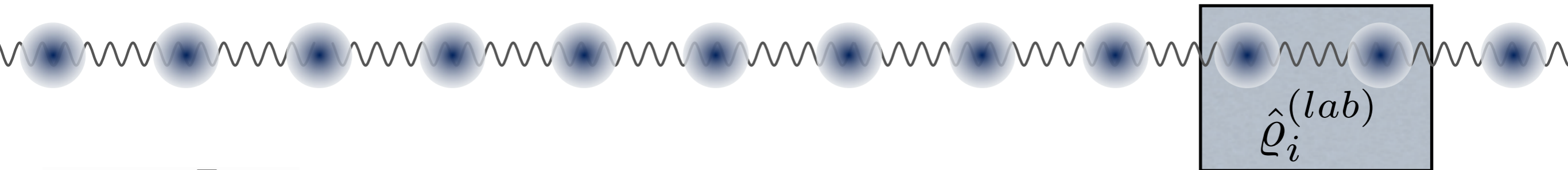
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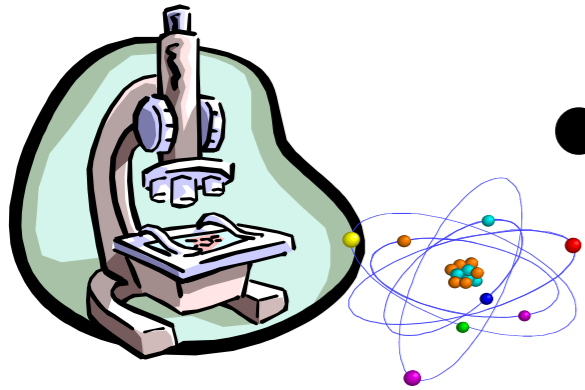
- Number and accuracy of measurements



Take only $4^2 N$



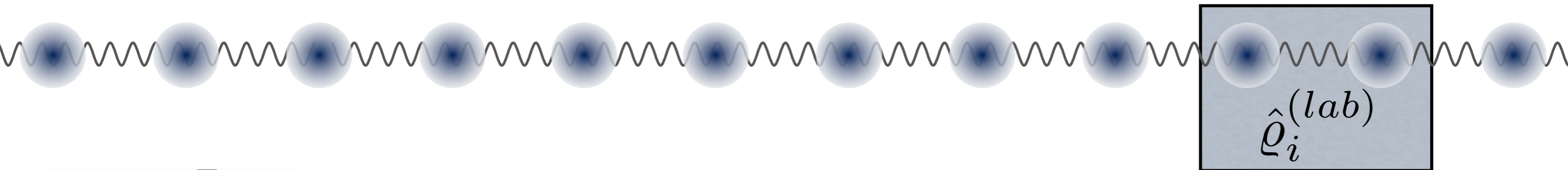
- *Promise:* $\hat{Q}_{lab} = |\psi\rangle\langle\psi|$ is unique ground state of local Hamiltonian $\hat{H} = \sum_i \hat{h}_i$:



- Number and accuracy of measurements



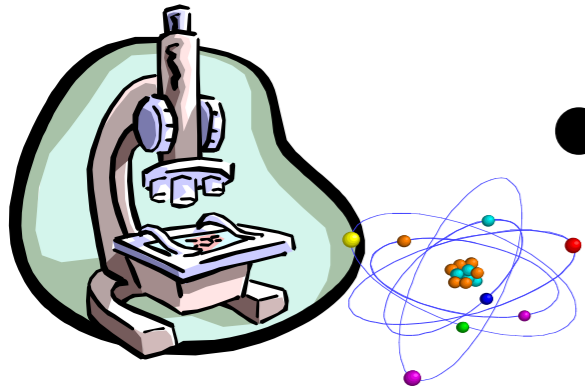
Take only $4^2 N$



- *Promise:* $\hat{Q}_{lab} = |\psi\rangle\langle\psi|$ is unique ground state of local Hamiltonian $\hat{H} = \sum_i \hat{h}_i$:

Candidate $\hat{Q}_{cand} = |\phi\rangle\langle\phi|$ with

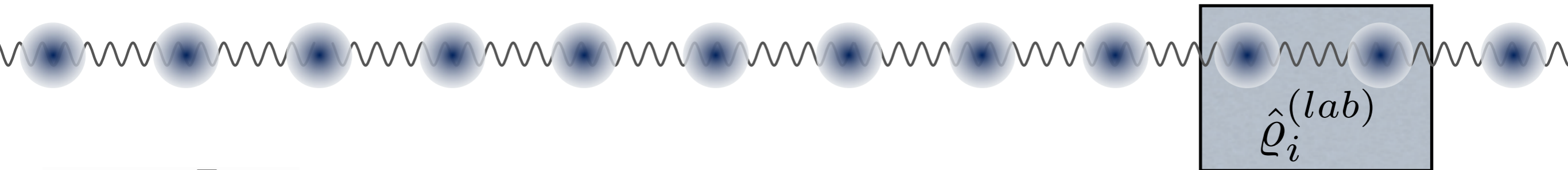
$$\hat{Q}_i^{(lab)} = \hat{Q}_i^{(cand)}$$



- Number and accuracy of measurements



Take only $4^2 N$

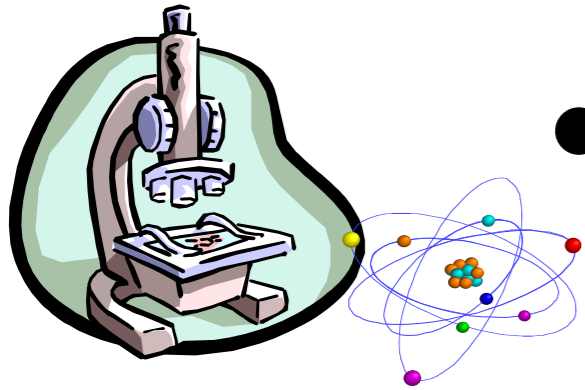


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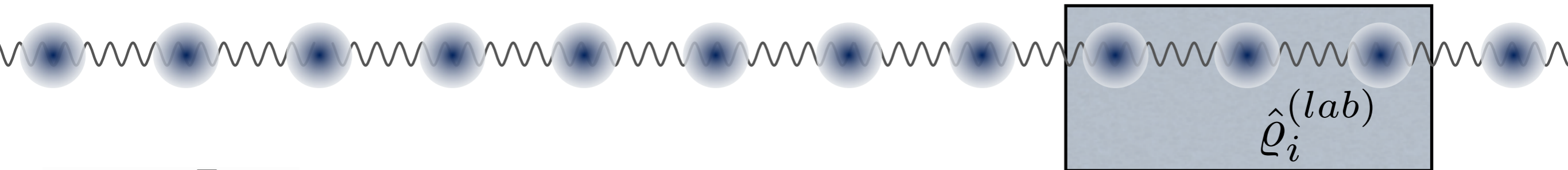
$$\longrightarrow \langle\psi|\hat{H}|\psi\rangle = \langle\phi|\hat{H}|\phi\rangle \longrightarrow |\psi\rangle = |\phi\rangle$$



- Number and accuracy of measurements



Take only $4^3 N$



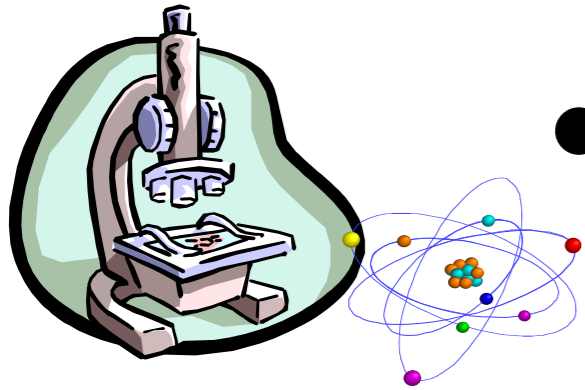
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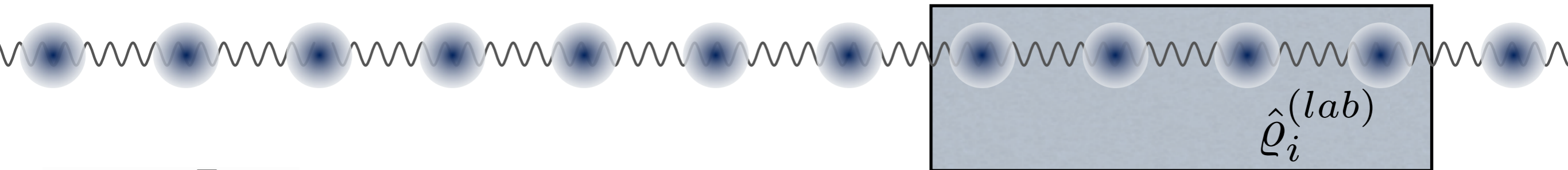
- General local Hamiltonians, limit



- Number and accuracy of measurements



Take only $4^4 N$



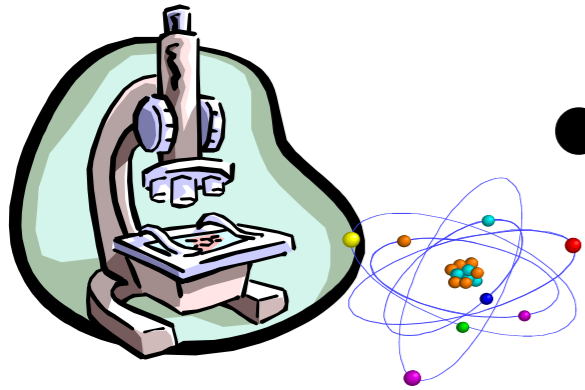
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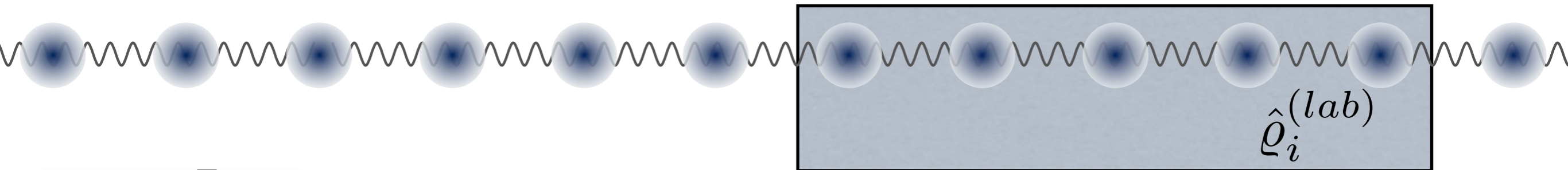
$$\hat{Q}_i^{(lab)} = \hat{Q}_i^{(cand)} \longrightarrow \langle\psi|\hat{h}_i|\psi\rangle = \langle\phi|\hat{h}_i|\phi\rangle$$

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- General local Hamiltonians, limit



- Number and accuracy of measurements \longrightarrow Take only $\sim N$



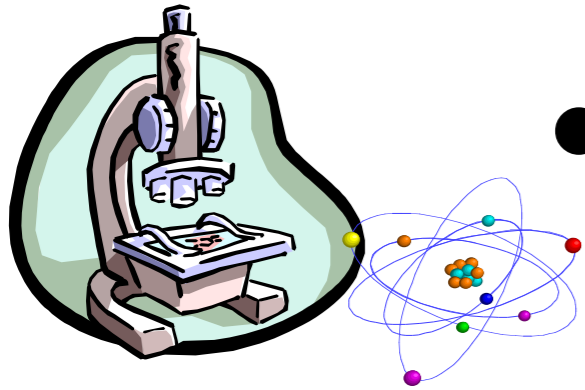
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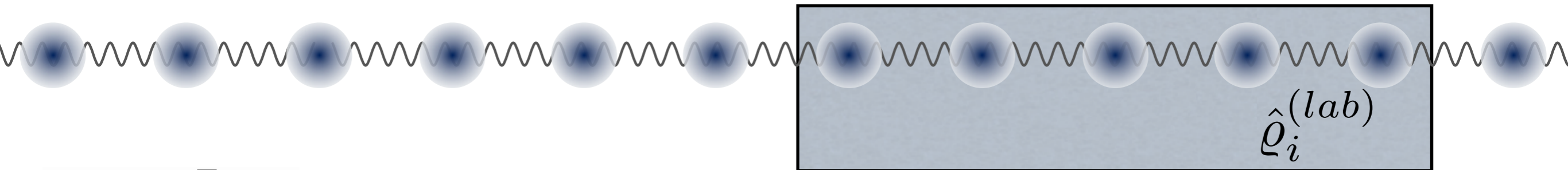
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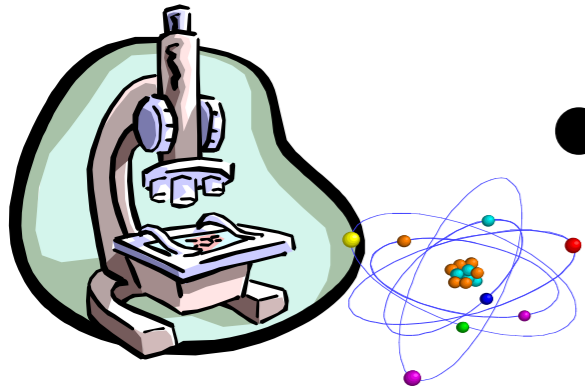
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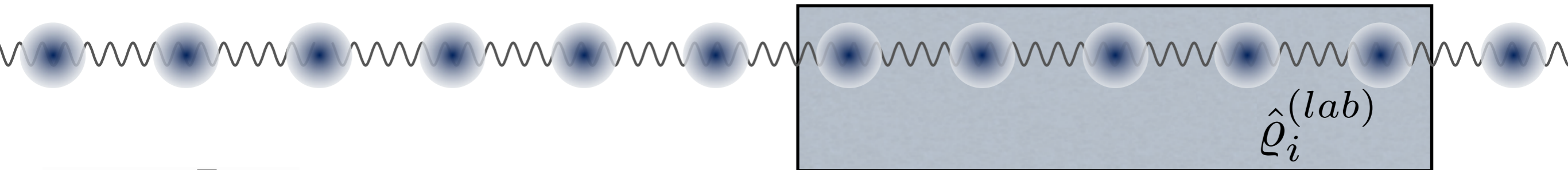
- Experiment: $\|\hat{Q}_i^{(lab)} - \hat{Q}_i^{(est)}\| \leq \epsilon_i$



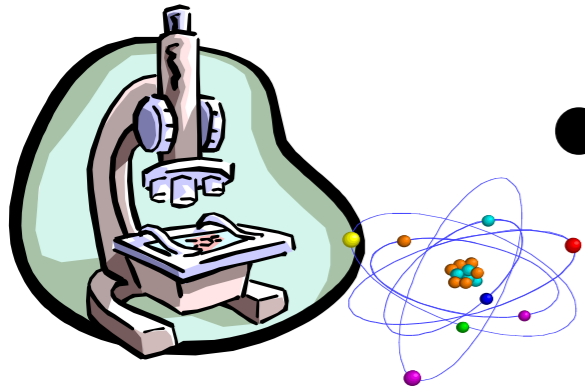
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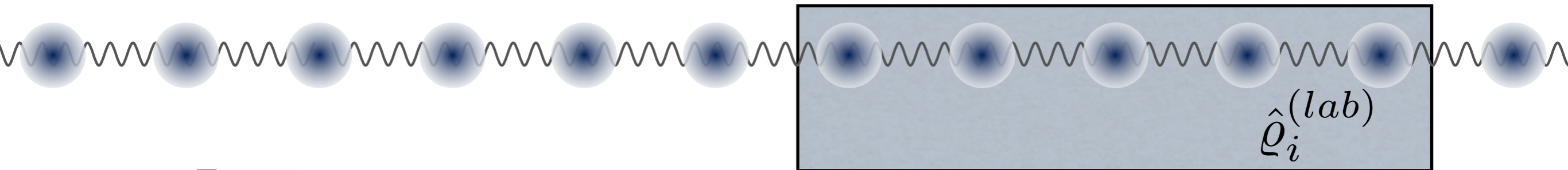
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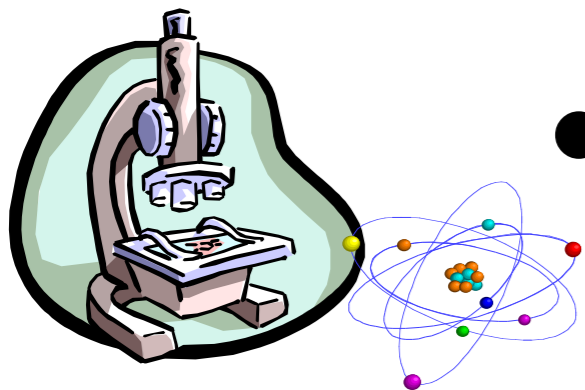
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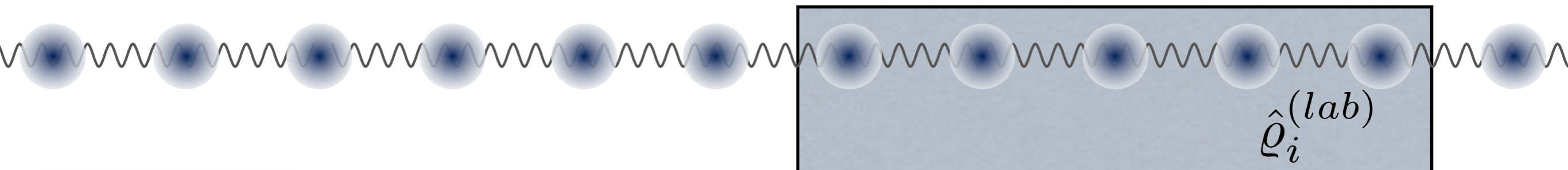
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- Set up *parent Hamiltonian* via $\hat{h}_i = \sum_{\lambda_k^{(i)}=0} |k\rangle_i \langle k|_i$
- $|\psi\rangle$ is ground state of $\hat{H} = \sum_i \hat{h}_i$



- Number and accuracy of measurements



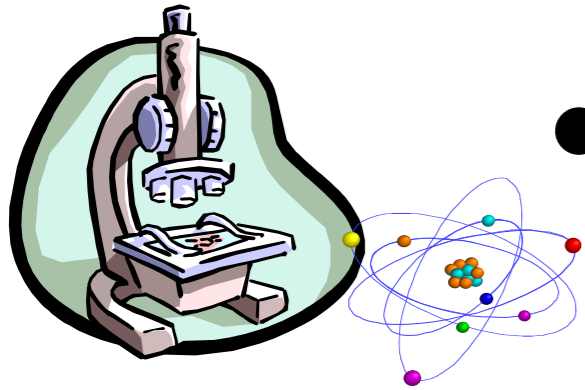
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If it is unique, we have

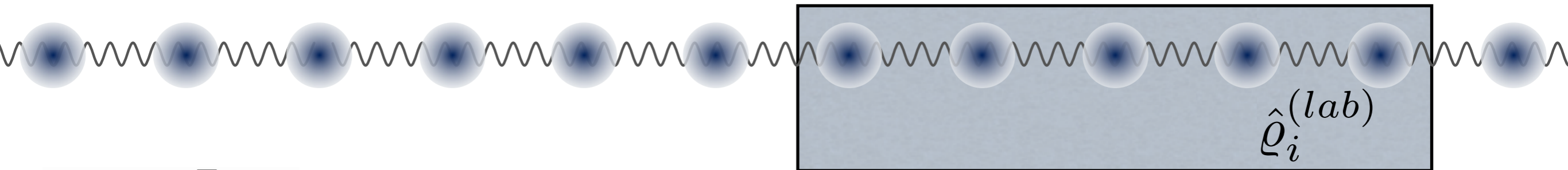
$$\text{tr}[\hat{H} \hat{Q}_{lab}] = \sum_{E_n > 0} E_n \langle n | \hat{Q}_{lab} | n \rangle \geq \Delta E \sum_{E_n > 0} \langle n | \hat{Q}_{lab} | n \rangle$$



- Number and accuracy of measurements



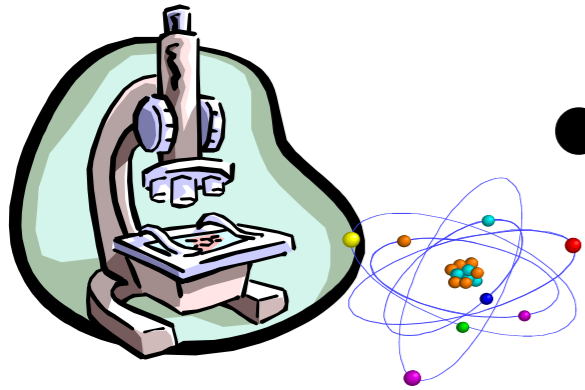
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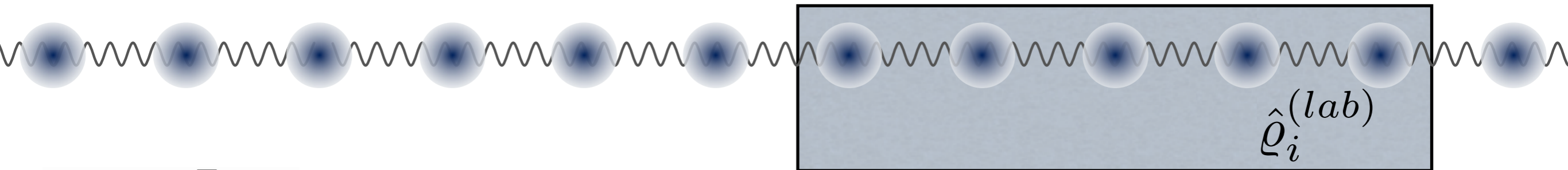
$$\begin{aligned} \text{tr}[\hat{H} \hat{Q}_{lab}] &= \sum_{E_n > 0} E_n \langle n | \hat{Q}_{lab} | n \rangle \geq \Delta E \sum_{E_n > 0} \langle n | \hat{Q}_{lab} | n \rangle \\ &= \Delta E (1 - \langle \psi | \hat{Q}_{lab} | \psi \rangle) \end{aligned}$$



- Number and accuracy of measurements



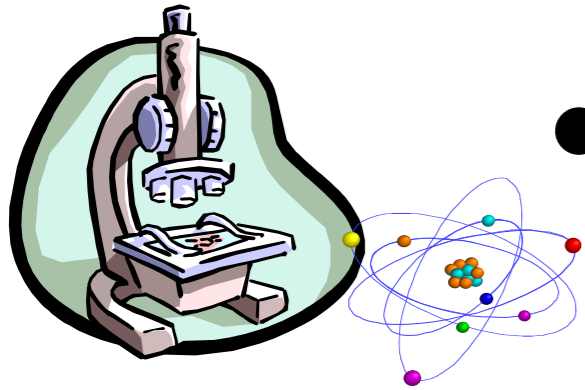
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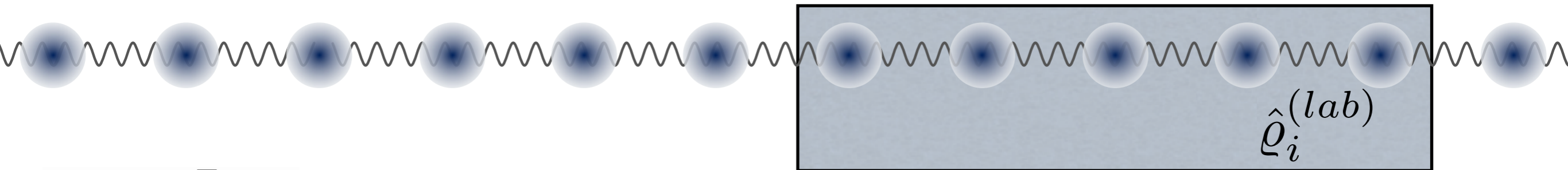
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$$1 - \langle\psi|\hat{Q}_{lab}|\psi\rangle \leq \frac{\text{tr}[\hat{H}\hat{Q}_{lab}]}{\Delta E}$$



- Number and accuracy of measurements

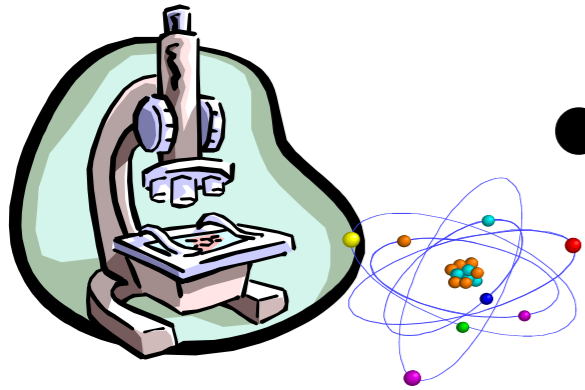
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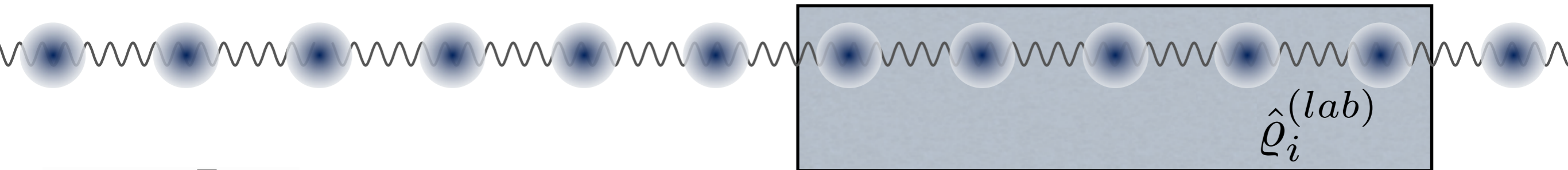
If it is unique, we have

$$1 - \langle\psi|\hat{Q}_{lab}|\psi\rangle \leq \frac{\text{tr}[\hat{H}\hat{Q}_{lab}]}{\Delta E} = \frac{\sum_i \text{tr}[\hat{h}_i(\hat{Q}_i^{(lab)} - \hat{Q}_i^{(est)})] + \text{tr}[\hat{h}_i\hat{Q}_i^{(est)}]}{\Delta E}$$



- Number and accuracy of measurements

→ Take only $\sim N$

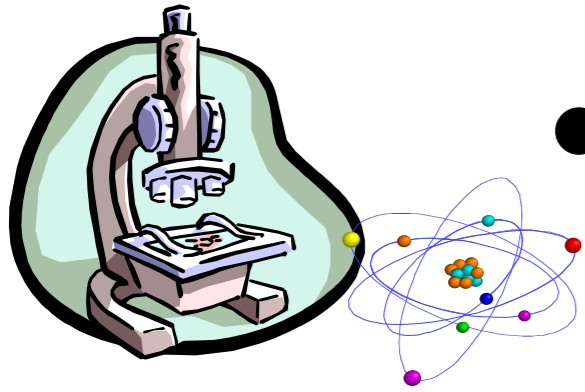


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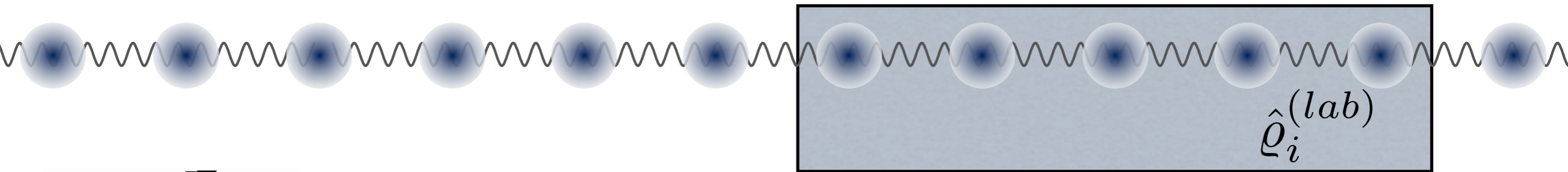
$$\leq \frac{\sum_i (\epsilon_i + \text{tr}[\hat{h}_i \hat{Q}_i^{(est)}])}{\Delta E}$$



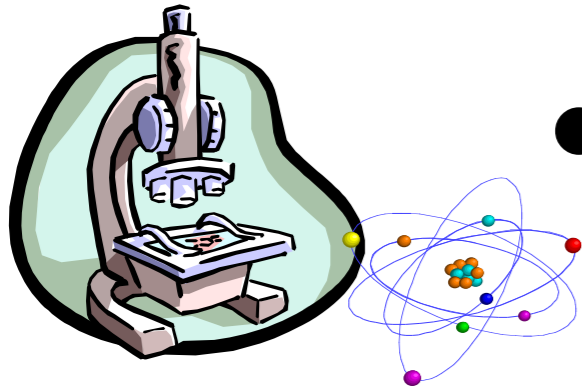
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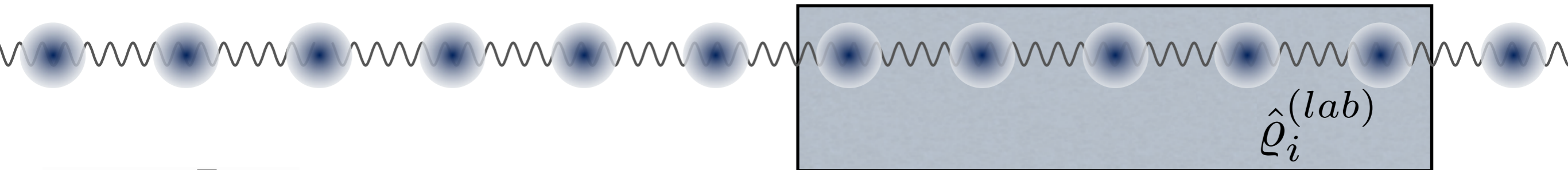
- Find compatible state
- Storage space



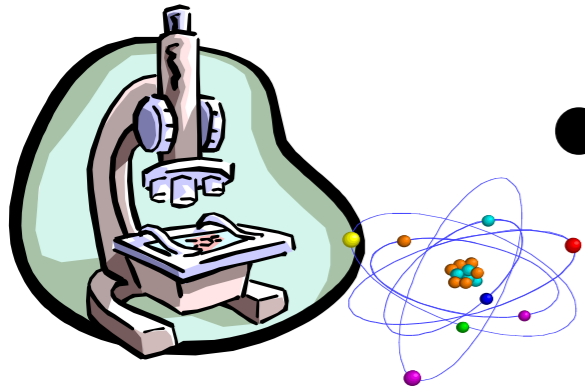
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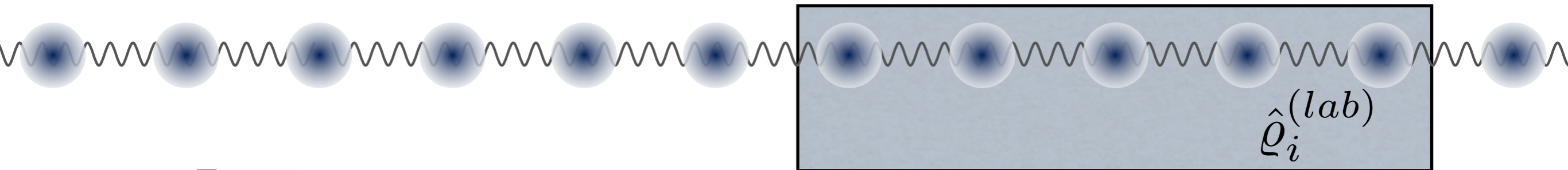
- Find compatible state
- Storage space
- Find parent Hamiltonian, ensure uniqueness of ground state and compute gap efficiently



- Number and accuracy of measurements

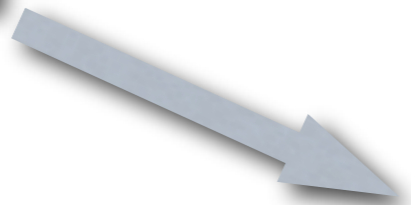


Take only $\sim N$



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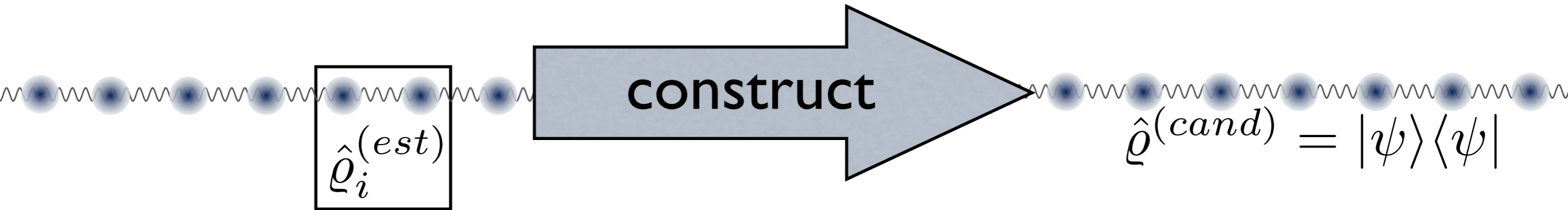
$$\|\hat{Q}_i^{(lab)} - \hat{Q}_i^{(est)}\| \leq \epsilon_i$$



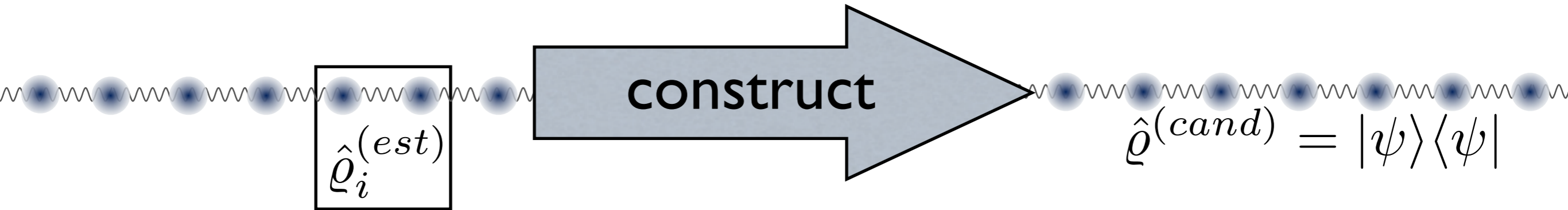
$$|\psi\rangle \text{ and } 1 - \langle \psi | \hat{Q}_{lab} | \psi \rangle \leq \frac{\sum_i (\epsilon_i + \text{tr}[\hat{h}_i \hat{Q}_i^{(est)}])}{\Delta E}$$

- Find compatible state

- Find compatible state, i.e., $|\psi\rangle$ such that $\hat{\rho}_i^{(est)} = \hat{\rho}_i^{(cand)}$



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measured Entries:

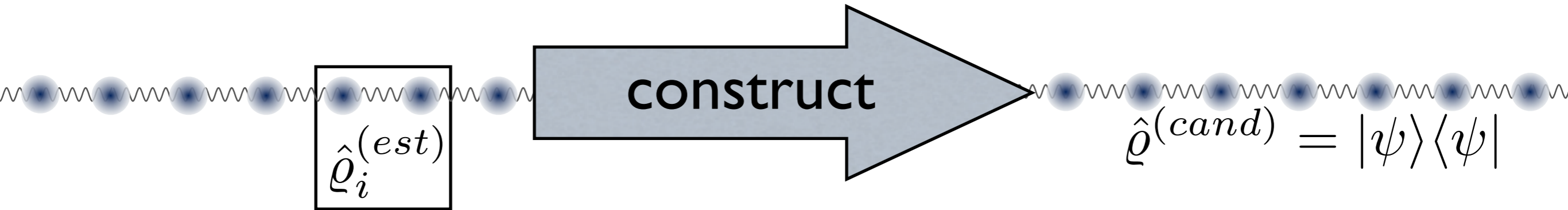
$$\Omega = \{k : \hat{P}_k = \mathbb{1} \otimes \hat{\sigma}_i^{\alpha_i} \otimes \hat{\sigma}_{i+1}^{\alpha_{i+1}} \otimes \mathbb{1}\}$$

Entries:

$$p_k = \text{tr}[\hat{P}_k |\psi\rangle\langle\psi|]$$

$$\hat{P}_k = \bigotimes_{i=1}^N \hat{\sigma}_i^{\alpha_i}$$

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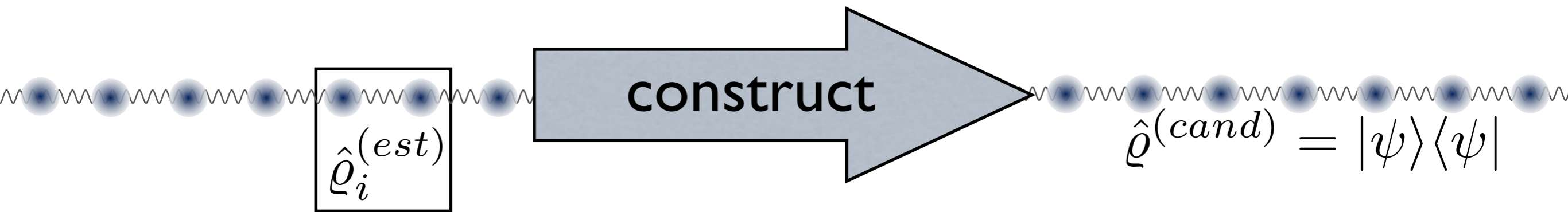
$$p_k = \text{tr}[\hat{P}_k |\psi\rangle\langle\psi|]$$

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Construct matrix from few known entries

- Find compatible state, i.e., $|\psi\rangle$ such that $\hat{\rho}_i^{(est)} = \hat{\rho}_i^{(cand)}$



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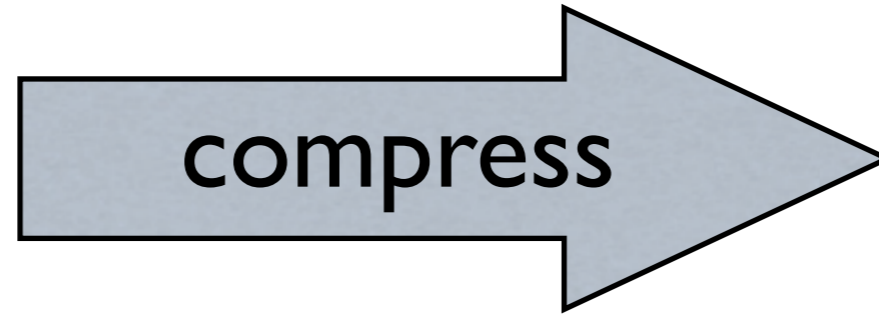
→ Construct matrix from few known entries

→ Compressed sensing

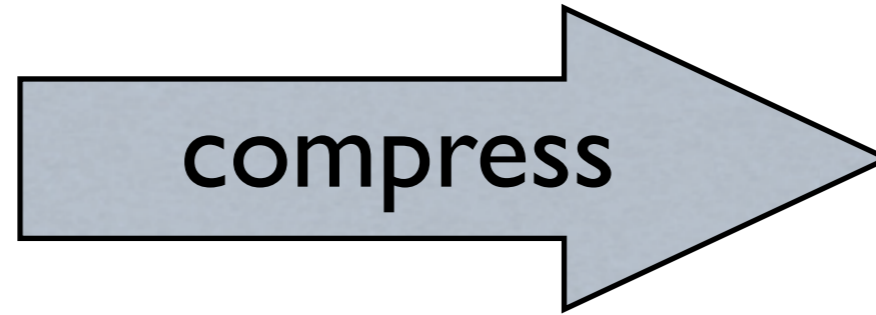
“a technique that may be the hottest topic in applied math today”

“paradigm-busting field in mathematics that’s reshaping the way people work with large data sets”

“Only six years old, compressed sensing has already inspired more than a thousand papers and pulled in millions of dollars in federal grants”



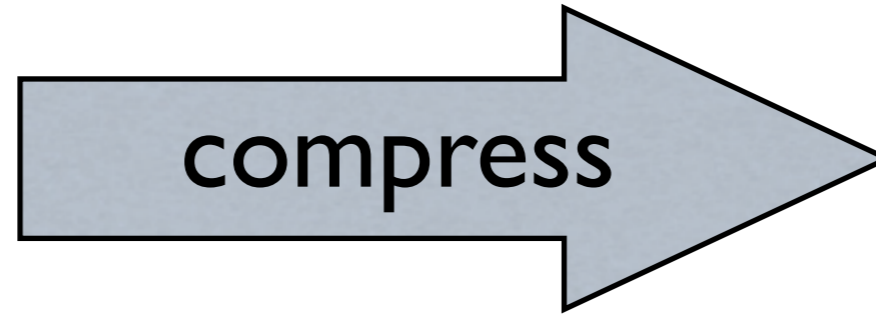
Pioneered by Donoho, Candès, Tao; an introduction: E.J. Candes and M.B. Wakin, *IEEE Sig. Proc. Mag.* **25**, 21 (2008); a quantum version: D. Gross, Y.-K. Liu, S.T. Flammia, S. Becker, J. Eisert, arXiv:0909.3304



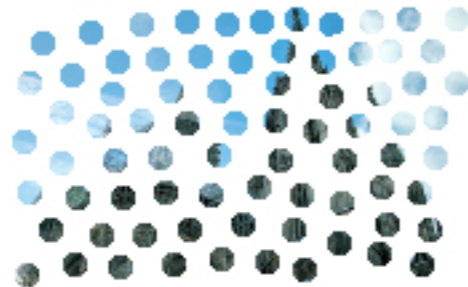
instead:



Pioneered by Donoho, Candès, Tao; an introduction: E.J. Candes and M.B. Wakin, *IEEE Sig. Proc. Mag.* **25**, 21 (2008); a quantum version: D. Gross, Y.-K. Liu, S.T. Flammia, S. Becker, J. Eisert, arXiv:0909.3304



instead:



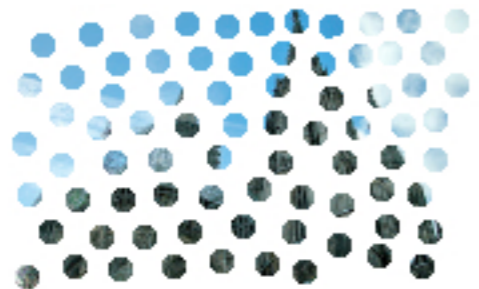
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compress

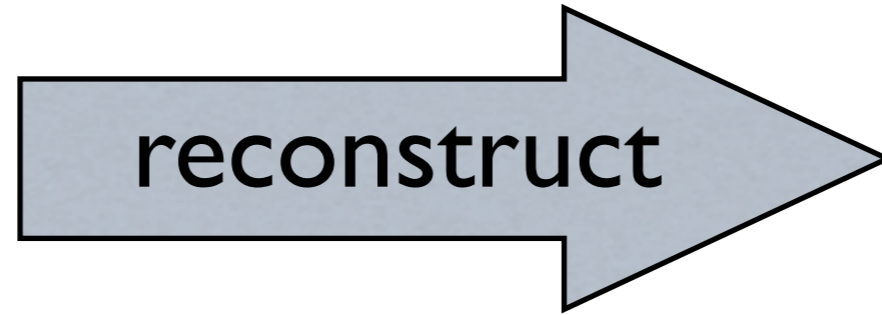
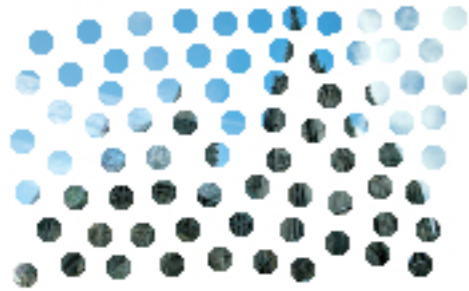


instead:



reconstruct





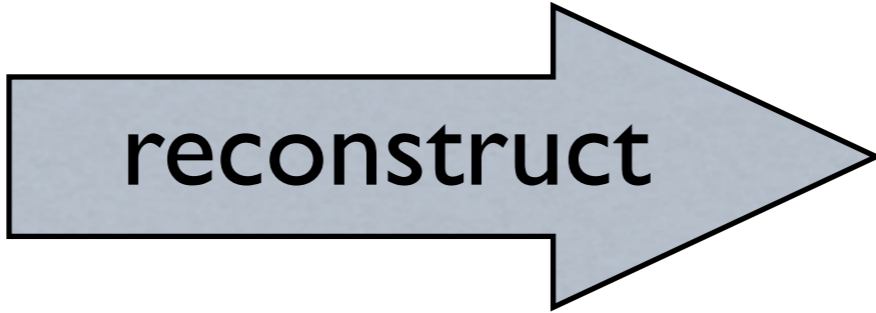
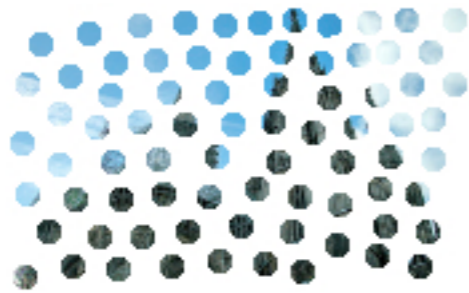
$m \times m$

sample size

$$\geq Crm^{1.2} \log(m)$$



small rank



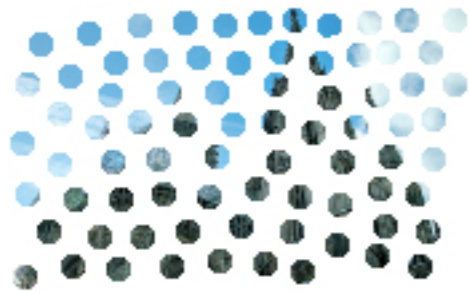
$m \times m$

sample size

$$\geq Crm^{1.2} \log(m)$$

small rank

perfect reconstruction
with very high probability



reconstruct

 $= M$ $m \times m$

sample size

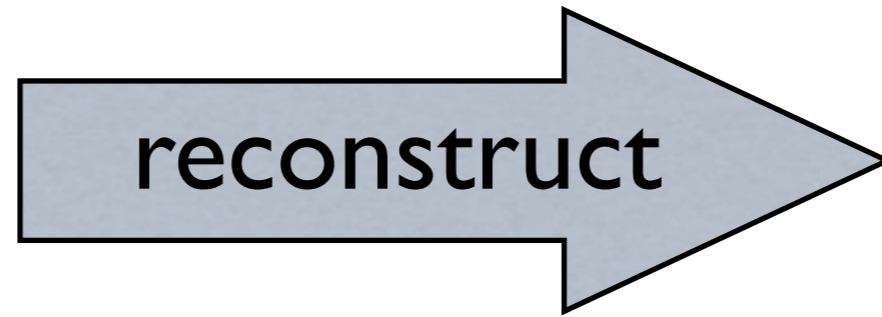
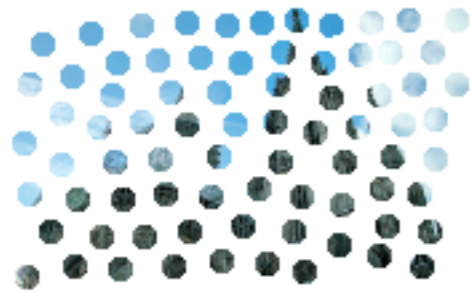
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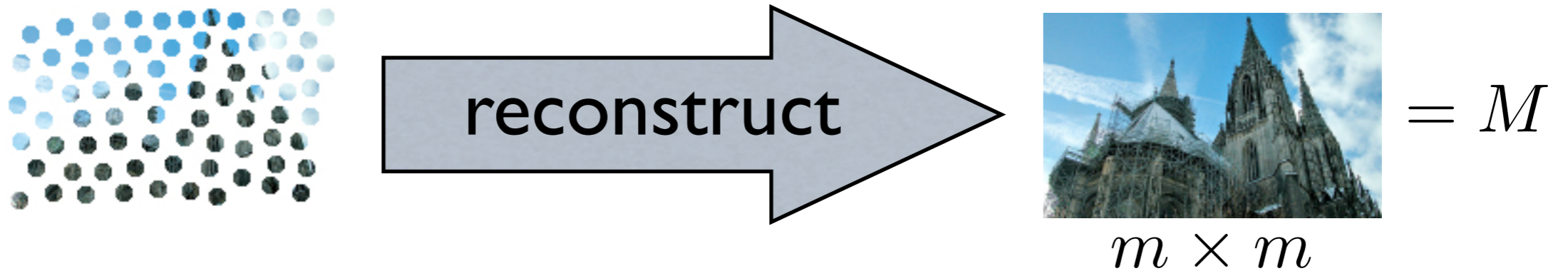
minimize $\text{tr}[|A|]$

such that $A_{i,j} = M_{i,j}$ for all $(i, j) \in \Omega$



$$= M$$

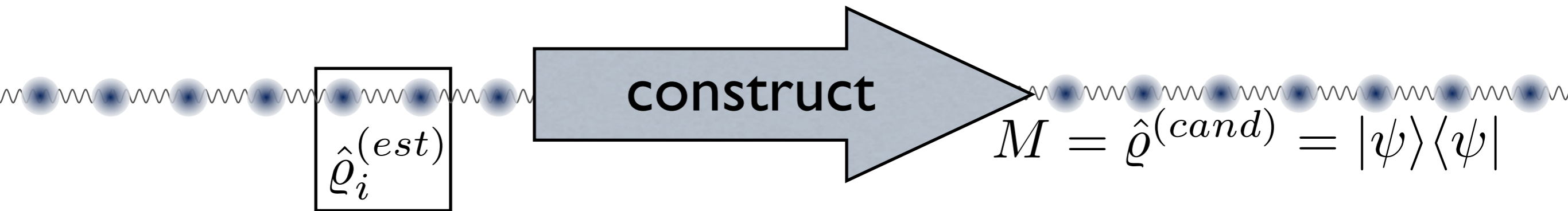
$$m \times m$$



Initialize Y_0 (e.g., by zero matrix), proceed inductively

- Singular value decomposition $Y_{n-1} = U\Sigma V^\dagger$
- Thresholding $X_n = U \max\{0, \Sigma - \mathbb{1}\tau\} V^\dagger$
- $Y_n = Y_{n-1} + \delta_n P_\Omega(M - X_n)$

Converges provably to solution for sufficiently small δ_n and $\tau \gg 1$



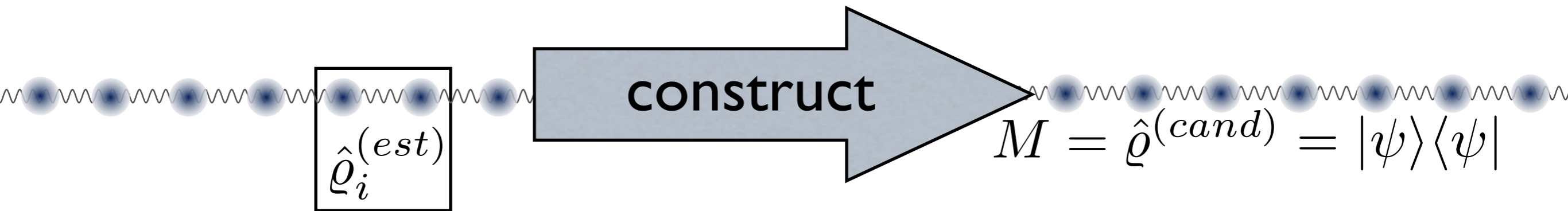
measured Entries:

$$\Omega = \{k : \hat{P}_k = \mathbb{1} \otimes \hat{\sigma}_i^{\alpha_i} \otimes \hat{\sigma}_{i+1}^{\alpha_{i+1}} \otimes \mathbb{1}\}$$

Entries:

$$p_k = \text{tr}[\hat{P}_k |\psi\rangle\langle\psi|]$$

$$\hat{P}_k = \bigotimes_{i=1}^N \hat{\sigma}_i^{\alpha_i}$$



measured Entries:

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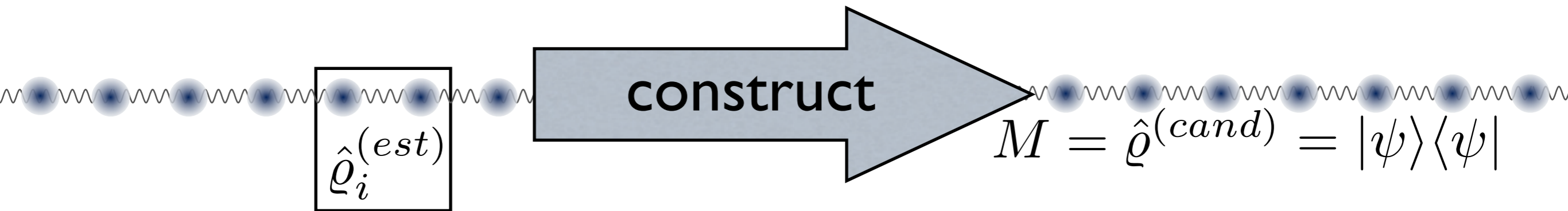
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Initialize Y_0 (e.g., by zero matrix), proceed inductively

- Singular value decomposition of $2^N \times 2^N$ matrix
- Thresholding $X_n = U \max\{0, \Sigma - \mathbb{1}\tau\} V^\dagger$
- $Y_n = Y_{n-1} + \delta_n \sum_{k \in \Omega} \frac{p_k - \text{tr}[X_n \hat{P}_k]}{2^N} \hat{P}_k$

Initialize Y_0 (e.g., by zero matrix), proceed inductively

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Instead of threshold, keep largest singular value

Initialize Y_0 (e.g., by zero matrix), proceed inductively

- **Thresholding** $|X_n\rangle = \|Y_{n-1}\| \cdot \operatorname{argmax} |\langle \phi | Y_{n-1} | \phi \rangle|$
- $Y_n = Y_{n-1} + \delta_n \sum_{k \in \Omega} \frac{p_k - \langle X_n | \hat{P}_k | X_n \rangle}{2^N} \hat{P}_k$

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local “Hamiltonian”

Initialize Y_0 (e.g., by zero matrix), proceed inductively

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local “Hamiltonian”, find ground state,

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local “Hamiltonian”, find ground state,
compute expectation values efficiently

DMRG, MPS methods

M. Fannes, B. Nachtergaele, and R.F.Werner, *Comm. Math. Phys.* **144**, 443 (1992),

U. Schollwöck, *Rev. Mod. Phys.* **77**, 259 (2005),

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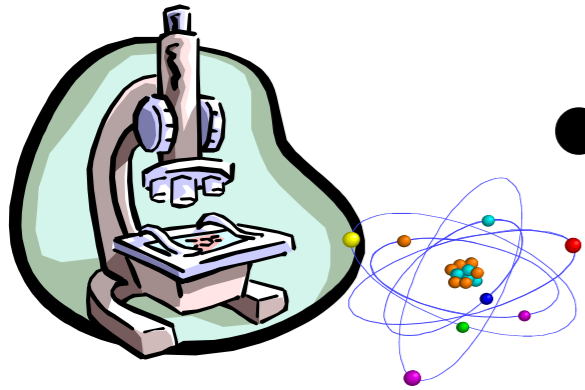
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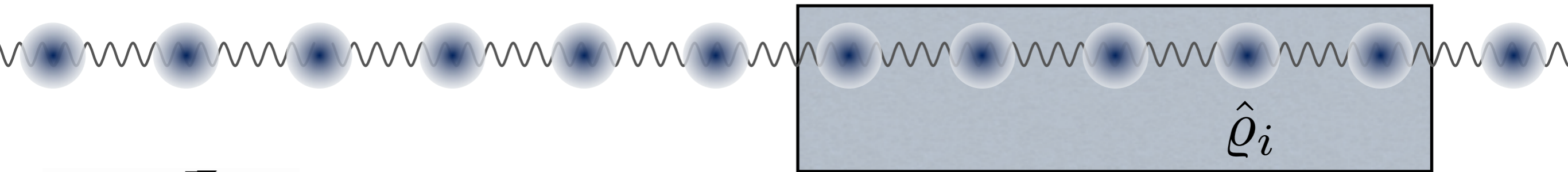
Matrix Product State $|MPS\rangle = \sum_{s_1, \dots, s_N} \operatorname{tr} [A_1[s_1] \cdots A_N[s_N]] |s_1 \cdots s_N\rangle$



- Number and accuracy of measurements



Take only $\sim N$

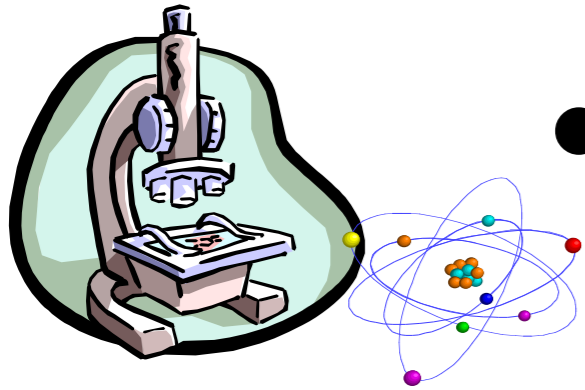


- Find compatible state
- Storage space
- Find *parent Hamiltonian*, ensure uniqueness of ground state and compute gap efficiently



MPS-SVT

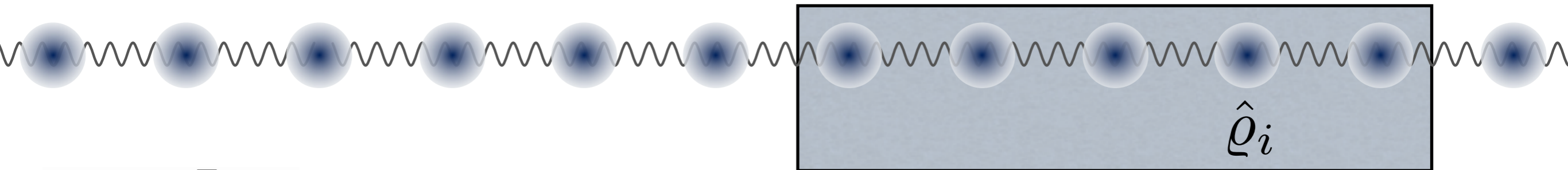
$2D^2 N$



- Number and accuracy of measurements



Take only $\sim N$



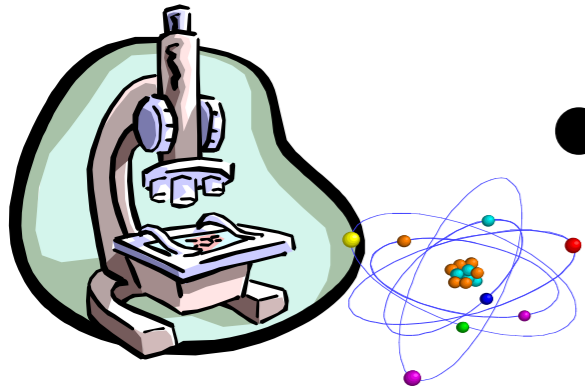
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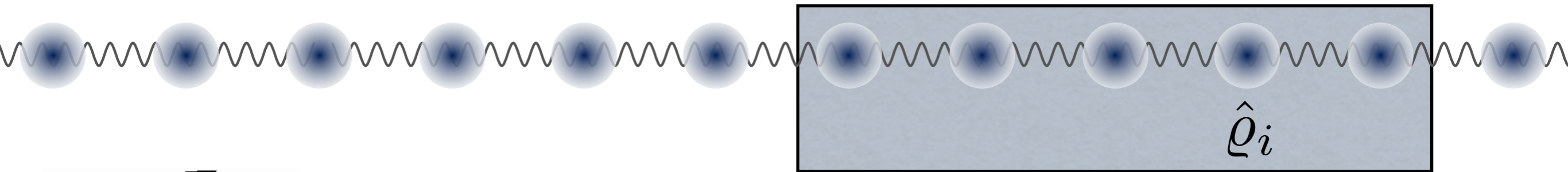




- Number and accuracy of measurements



Take only $\sim N$



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MPS-SVT

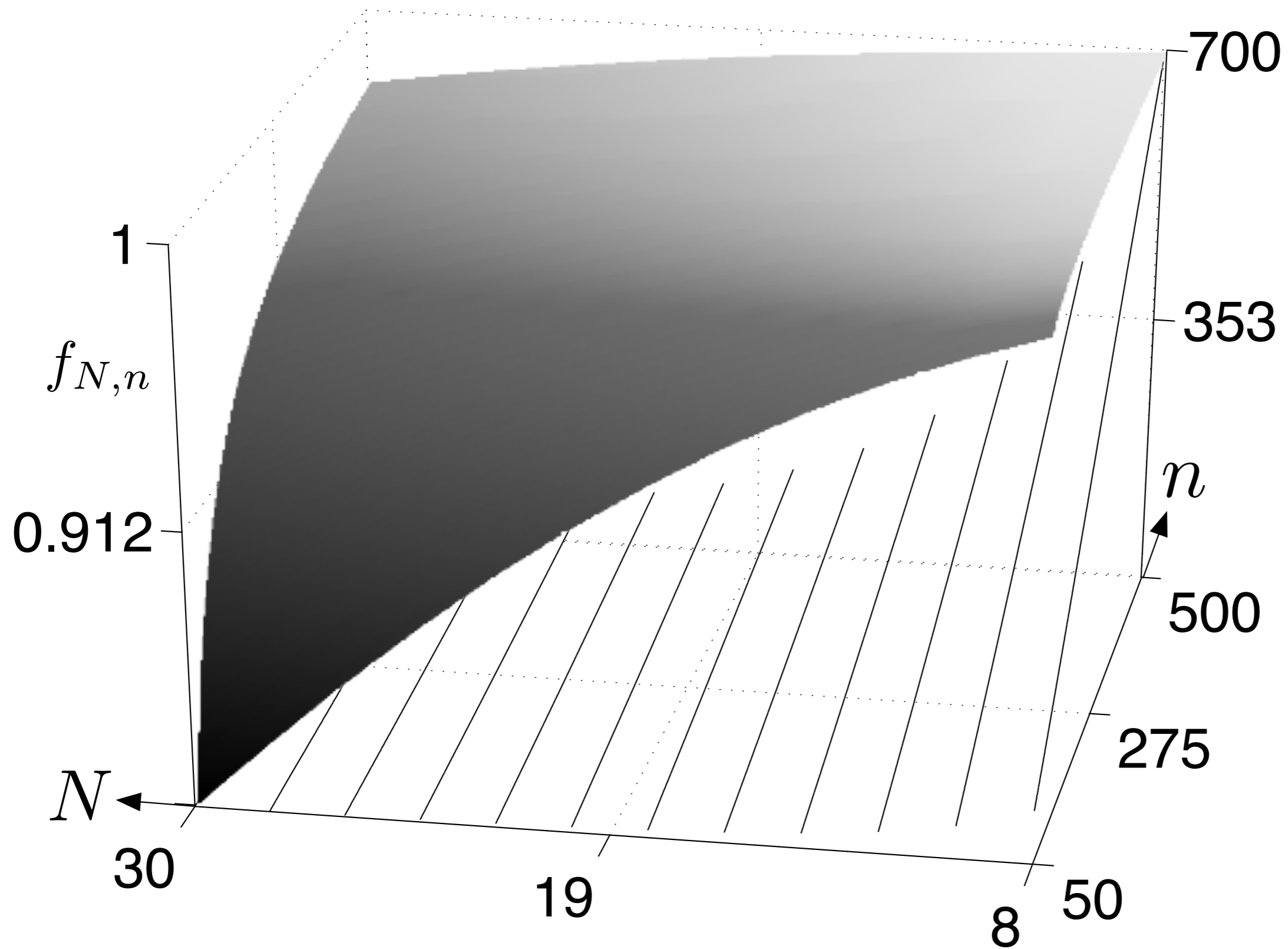
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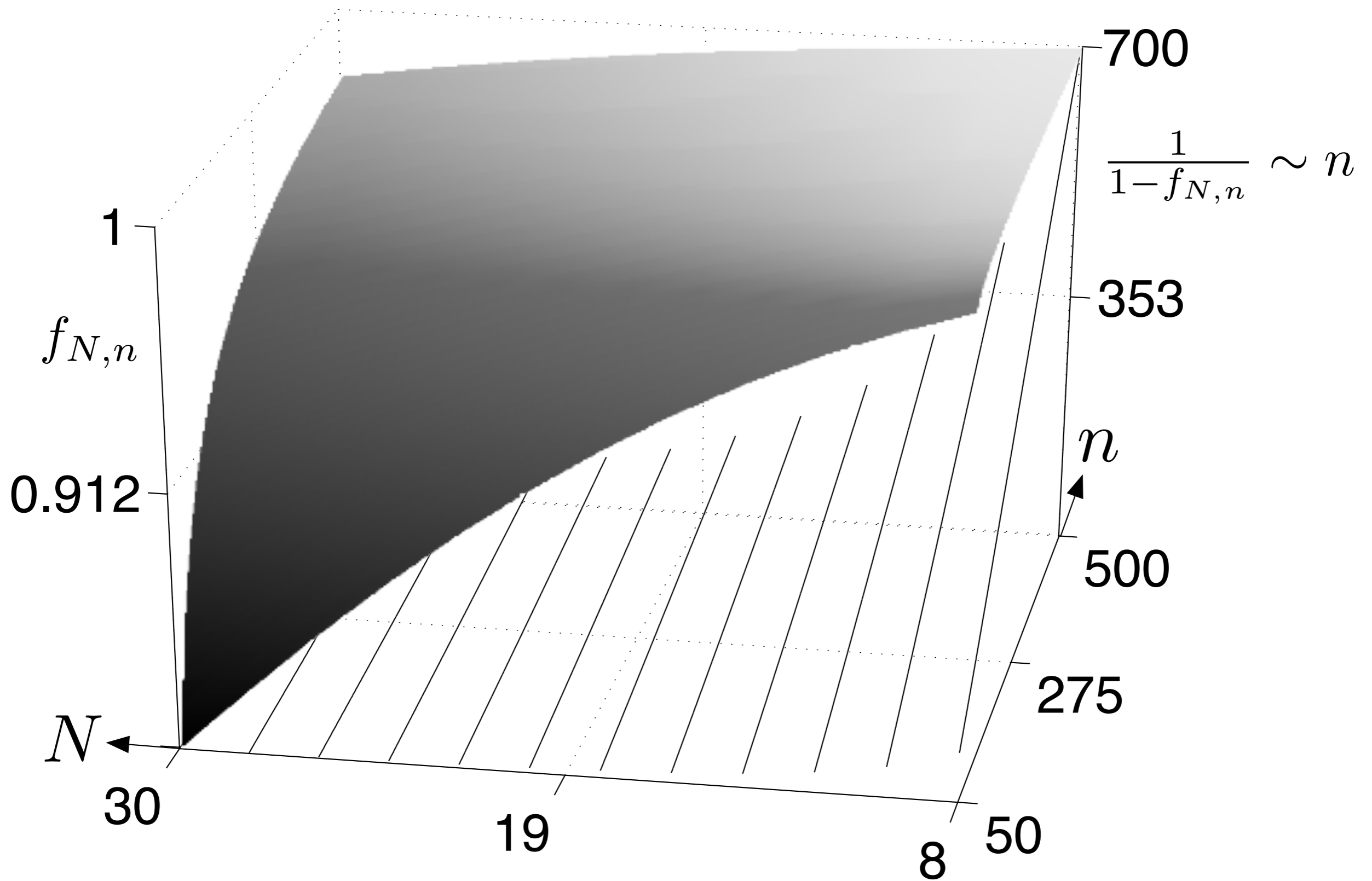


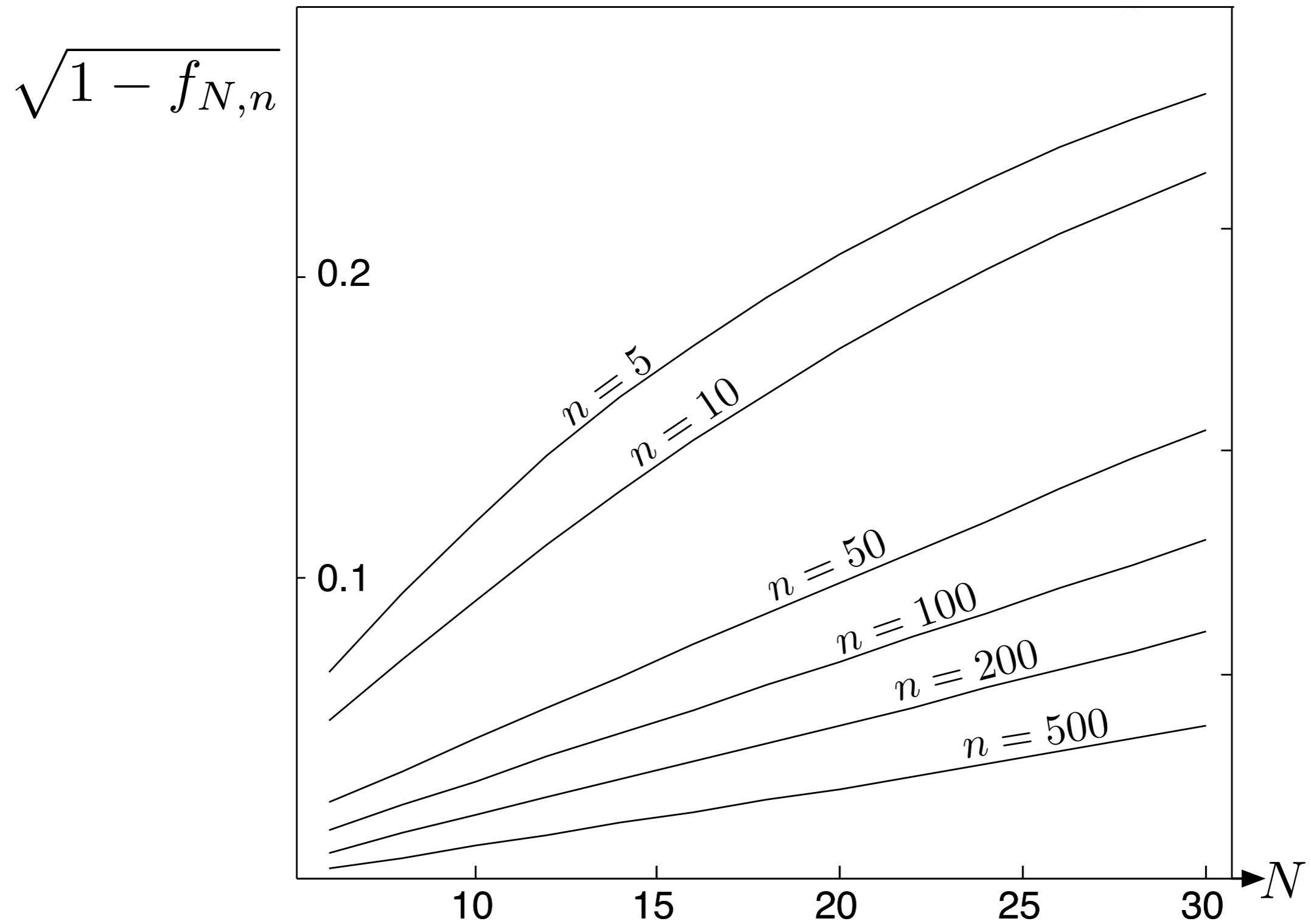
How does the algorithm perform?

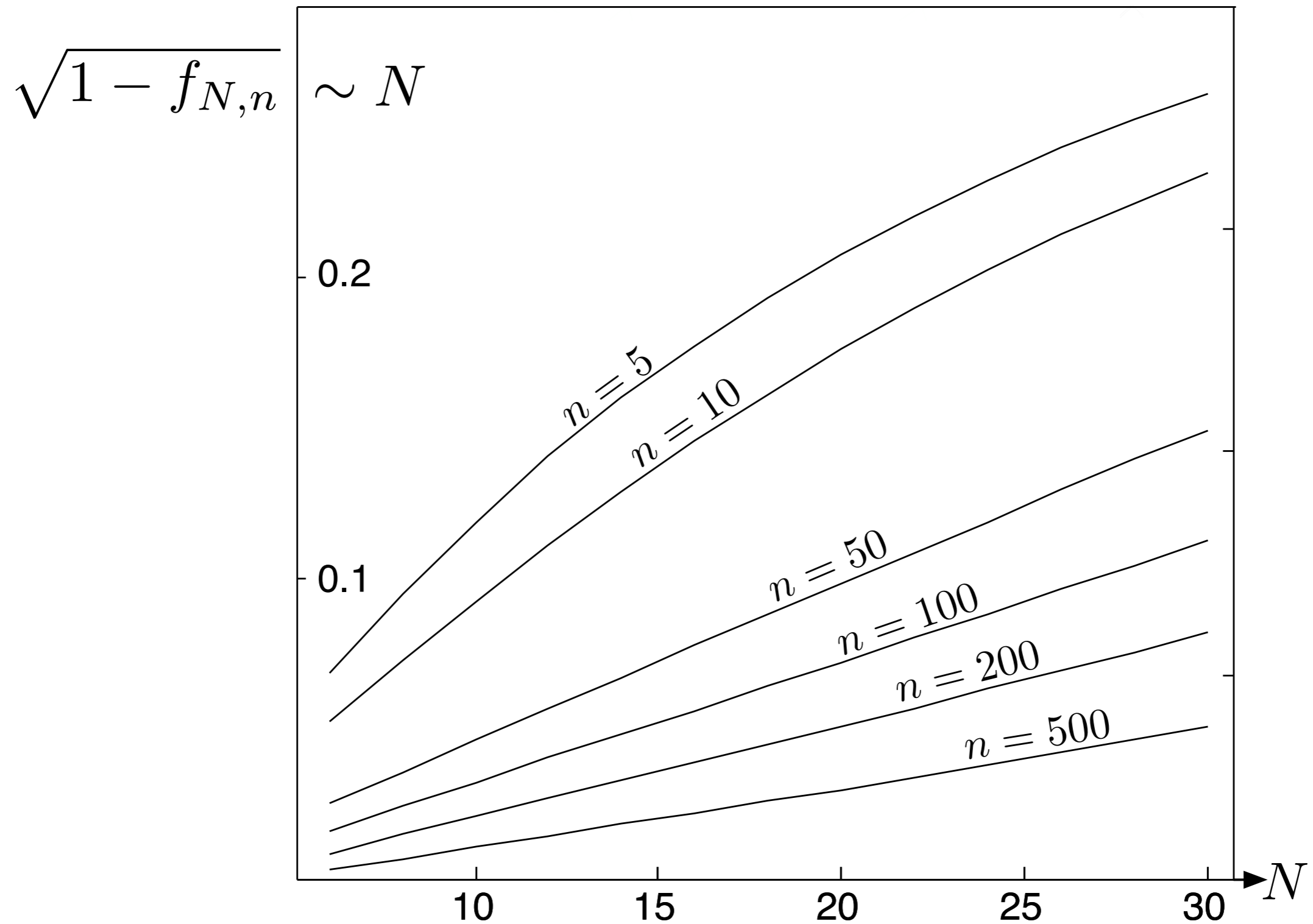
Ground state critical Ising model $\hat{H} = - \sum_{i=1}^{N-1} \hat{\sigma}_i^x \hat{\sigma}_{i+1}^x - \sum_{i=1}^N \hat{\sigma}_i^z$

- “Measure” all $\hat{Q}_{i,i+1}$
- Completely determines ground state
- Compute fidelity $f_{N,n} = |\langle gs | X_n \rangle|^2$







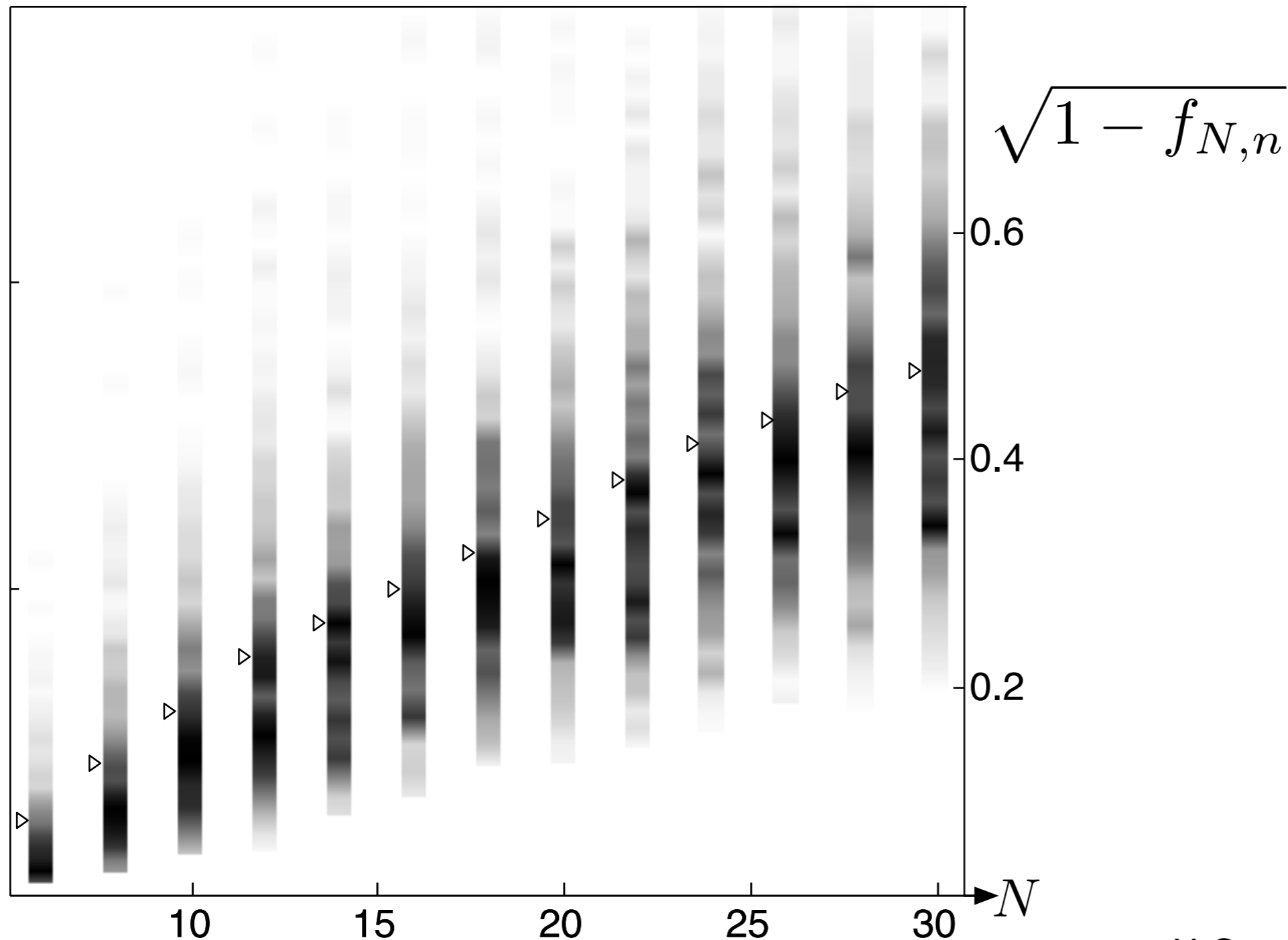


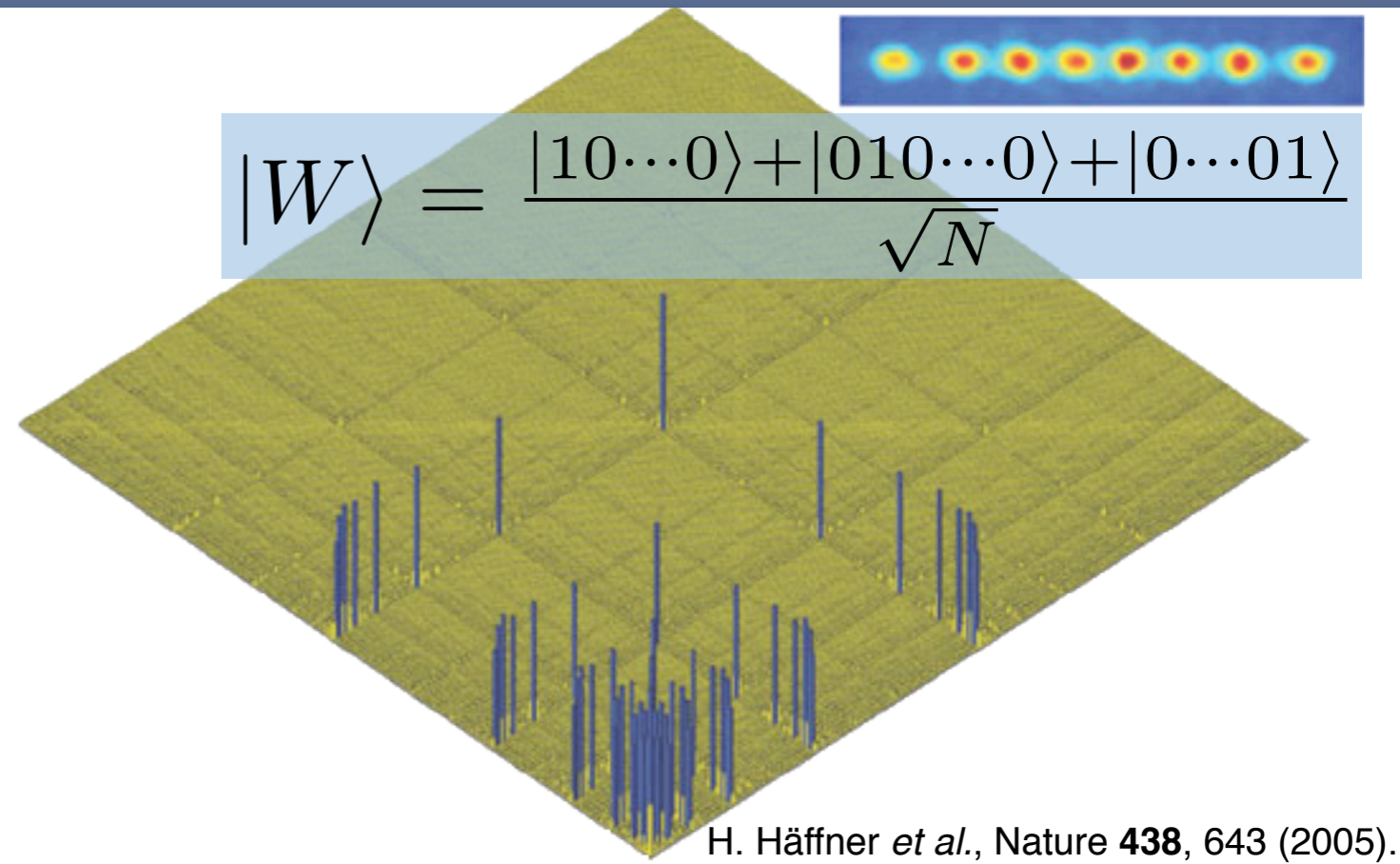
$$\hat{H} = \sum_{i=1}^{N-1} \hat{r}_i^{(i)} \hat{r}_{i+1}^{(i)}$$

hermitian, real and imaginary part of entries uniformly from $[-1, 1]$,
1000 realizations for each $N, n = 5$

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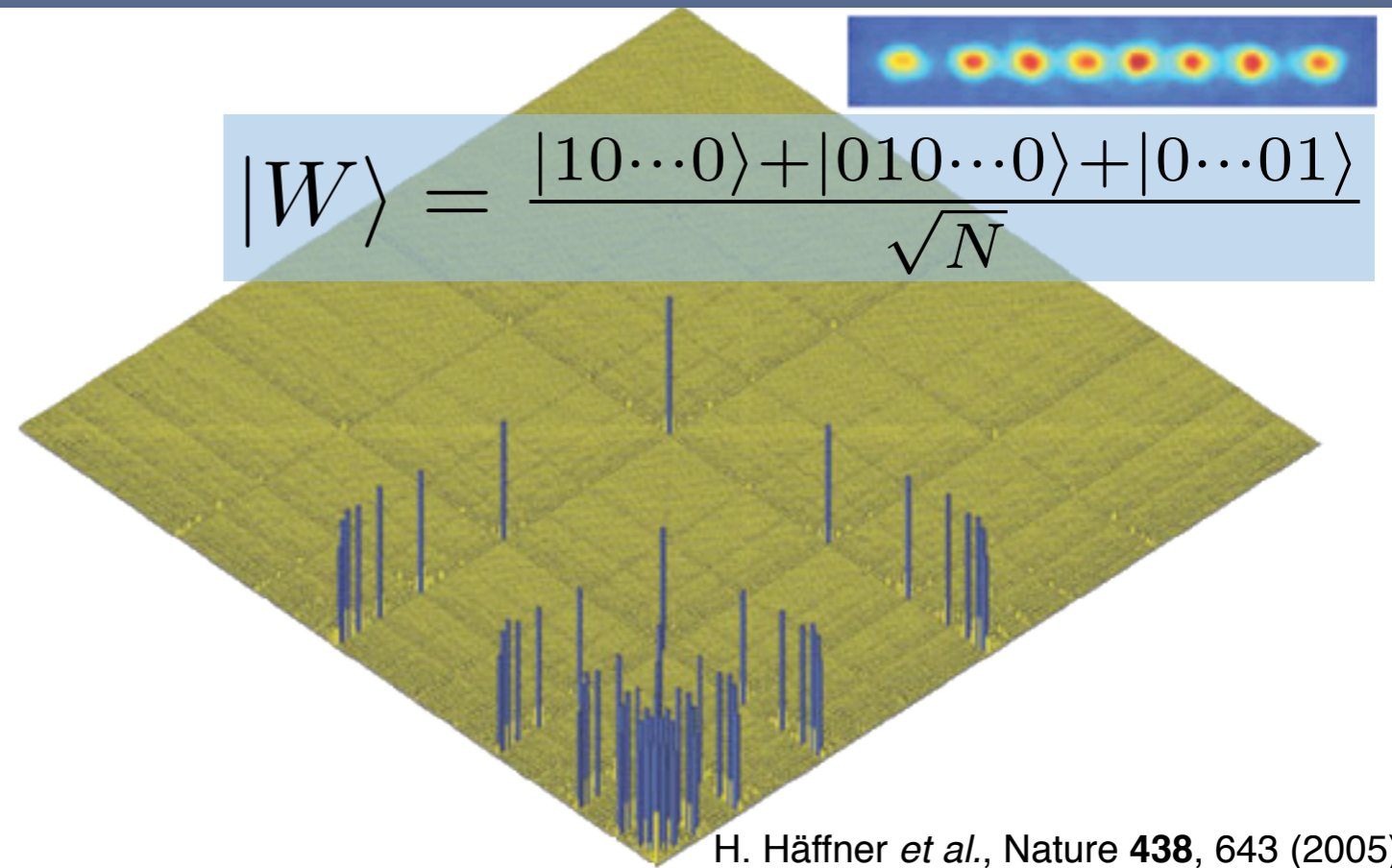




$$p_k = \langle W | \hat{P}_k | W \rangle + r$$

Gaussian, zero mean

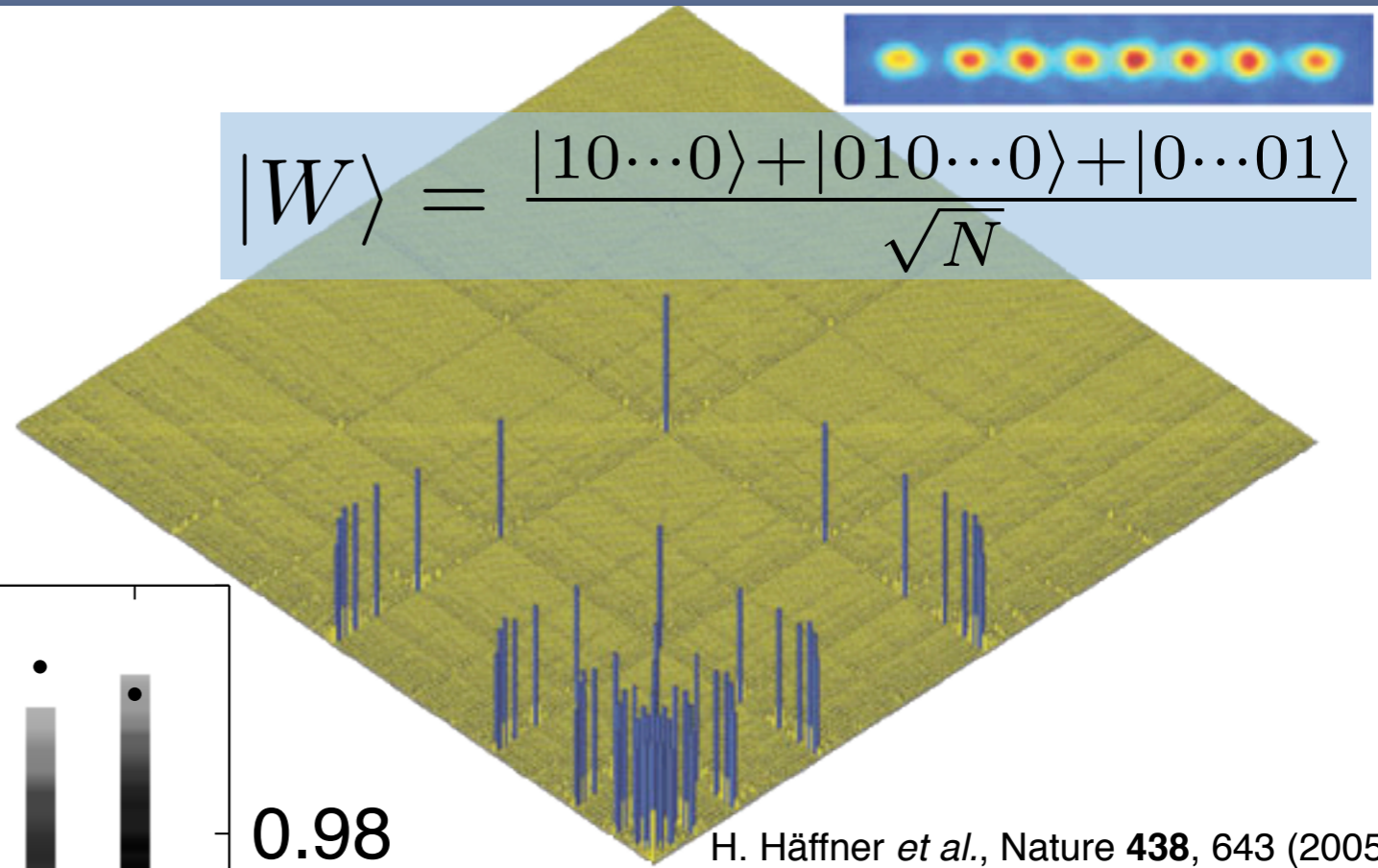
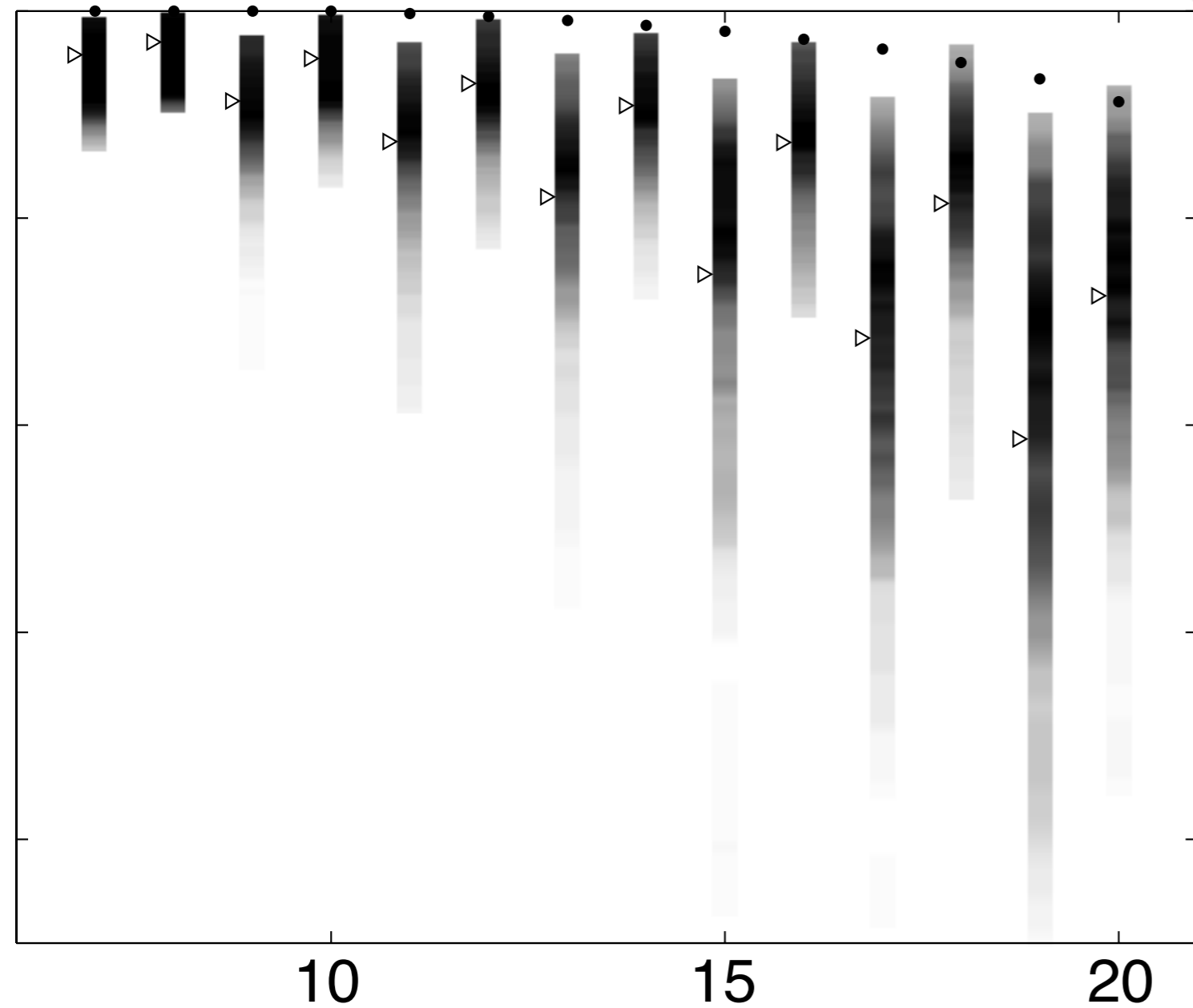
100 realizations for each N



$$p_k = \langle W | \hat{P}_k | W \rangle + r$$

Gaussian, zero mean

100 realizations for each N



0.98

0.96

0.94

0.92

N

$$f_{N,n} = |\langle W | X_n \rangle|^2$$

$n = 4000$

$\sigma = 0.01$ (odd N)

$\sigma = 0.005$ (even N)

- Efficient ($\text{poly}(N)$) scheme to reconstruct states from few measurements ($\sim N$)
- Extensive numerics suggest convergence. In fact, equipped with candidate MPS, fidelity can under certain conditions be bounded. S.T. Flammia, D. Gross, S.D. Bartlett, R. Somma, arXiv:1002.3839
- Generalization to higher dimensions: Analogues of MPS, e.g., MERA, PEPS,....
- Mixed states: Purification, e.g., Gibbs state minimizes $\text{tr}[\hat{\rho}\hat{H}] - TS(\hat{\rho})$ and entropy density $\lim_{k \rightarrow \infty} S(\hat{\rho}_i)/k$ exists